

TESTING THE MUON G-2 ANOMALY AT THE LHC

based on JHEP 1405(2014)145 with Ayres Freitas, Stefan Kell and Joe Lykken

Susanne Westhoff



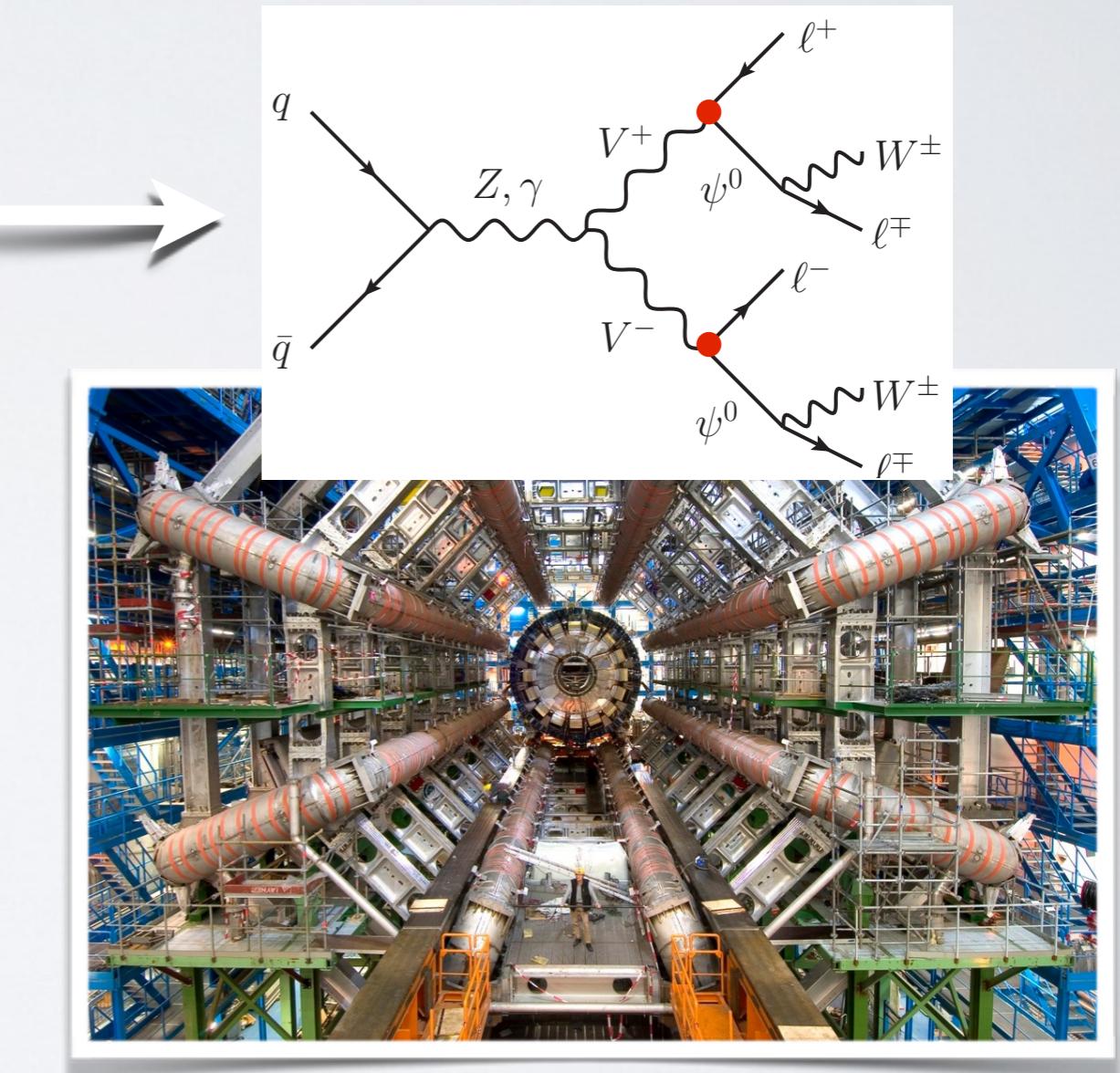
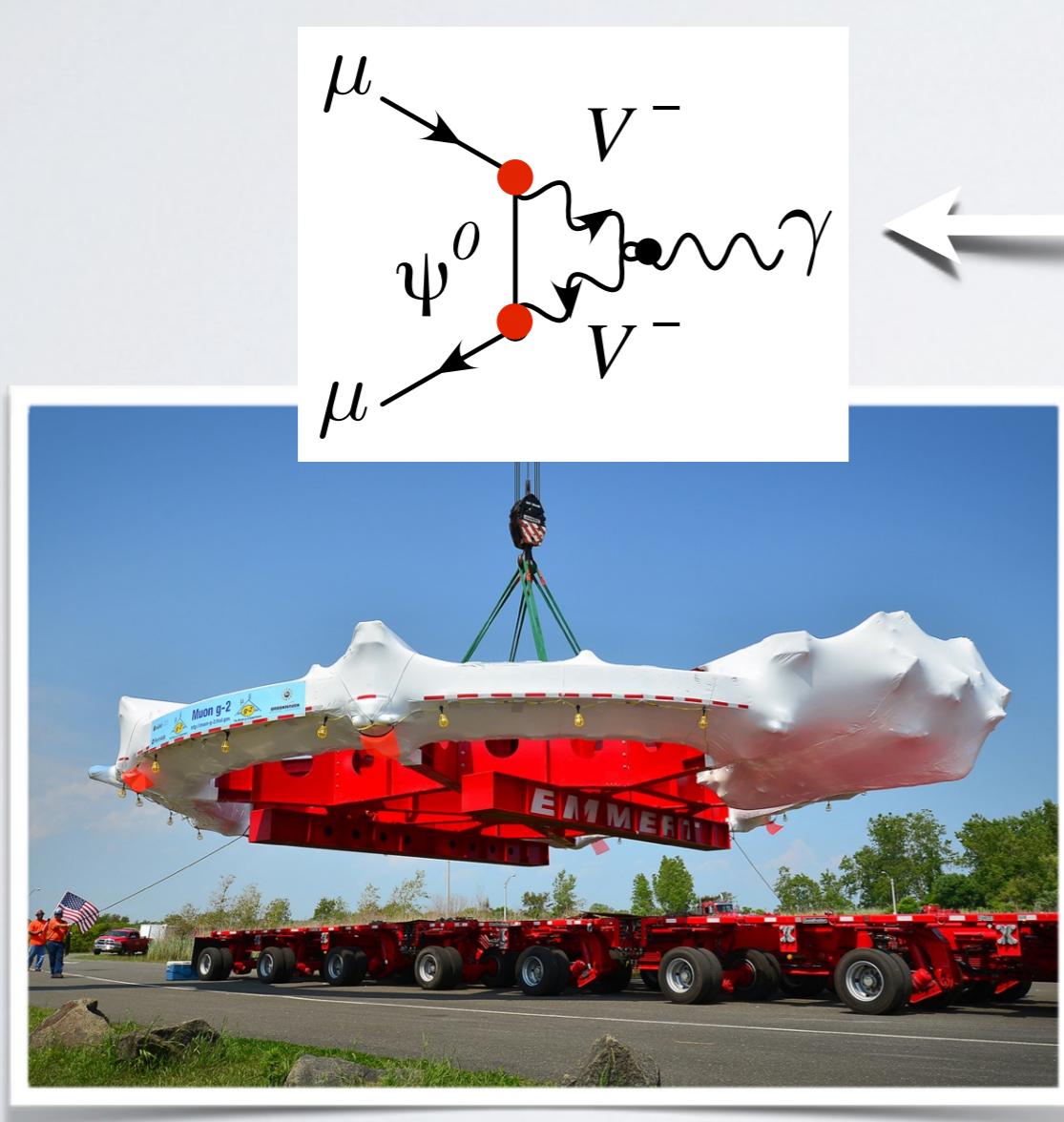
University of Pittsburgh

Theory Seminar - April 8, 2015 - Fermilab

MUON G-2 AND THE LHC

Lepton magnetic moment and its anomalous part:

$$\vec{\mu}_\ell = g_\ell \frac{e}{2m_\ell} \vec{s}, \quad a_\ell = \frac{g_\ell - 2}{2}$$



LEPTON ANOMALOUS MAGNETIC MOMENTS

Electron anomalous magnetic moment measured at Harvard:

$$a_e^{\text{exp}} = 1.15965218073(28) \times 10^{-3} \quad (\text{0.24 ppb})$$

[Hanneke, Fogwell, Gabrielse, PRL 100 (2008) 120801]

SM prediction including electromagn. corrections up to $(\alpha/\pi)^5$:

$$a_e^{\text{SM}} = 1.15965218178(77) \times 10^{-3}$$

[Aoyama et al., PRL 109 (2012) 111807]

$$\rightarrow \text{fine structure constant } \alpha^{-1} = 137.035999174(35)$$

Muon anomalous magnetic moment measured at Brookhaven:

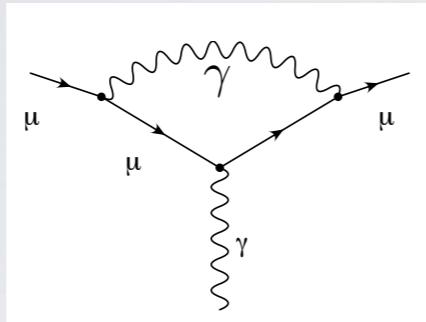
$$a_\mu^{\text{exp}} = 1.16592080(54)(33) \times 10^{-3} \quad (\text{0.54 ppm})$$

[BNL Muon g-2 collaboration, PRD73, 072003 (2006)]

MUON G-2 IN THE STANDARD MODEL

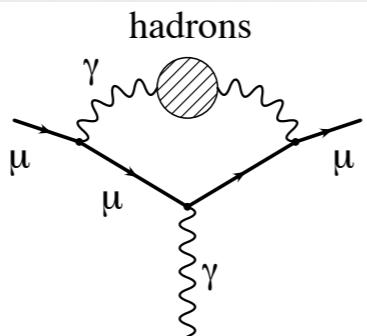
[Davier et al., Eur.Phys.J. C71(2011) 1515]

$$a_\mu^{\text{SM}} = a_\mu^{\text{QED}}$$



$$\frac{\alpha}{2\pi} + \dots \left(\frac{\alpha}{\pi}\right)^5 : 116584718.951(10) [10^{-11}]$$

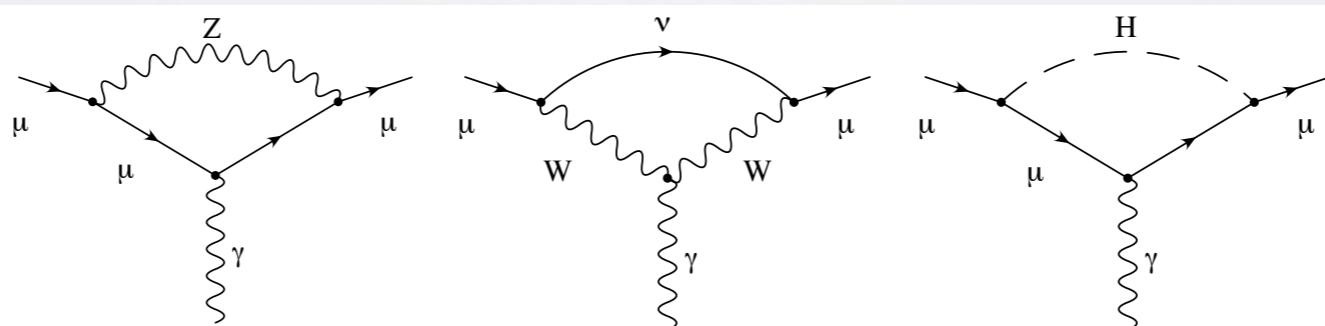
$$+ a_\mu^{\text{had}}$$



Leading order:
Higher orders:
Light-by-light sctg.:

$$6923(42) \\ -97.9(0.9) \\ 105(26)$$

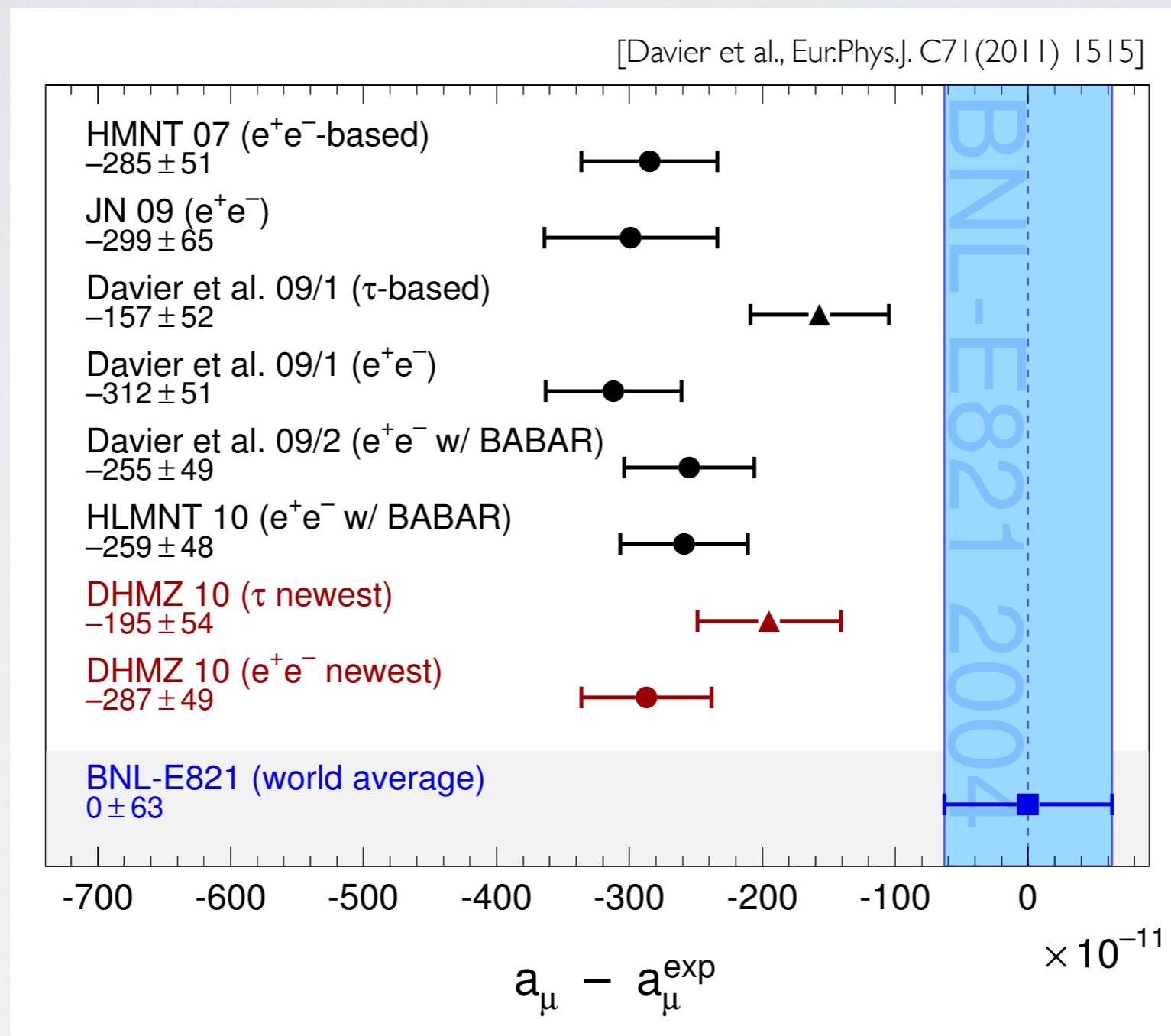
$$+ a_\mu^{\text{EW}}$$



$$\frac{G_F m_\mu^2}{8\sqrt{2}\pi^2} \left[\frac{5}{3} + \frac{1}{3}(1 - 4 \sin^2 \theta_W)^2 + \mathcal{O}\left(\frac{m_\mu^2}{M_{\text{EW}}^2}\right) \right] + (\text{2-loop}) : 154(1)$$

$$a_\mu^{\text{SM,tot}} = 116591802(49) \times 10^{-11}$$

CONFRONTED WITH THE MEASUREMENT



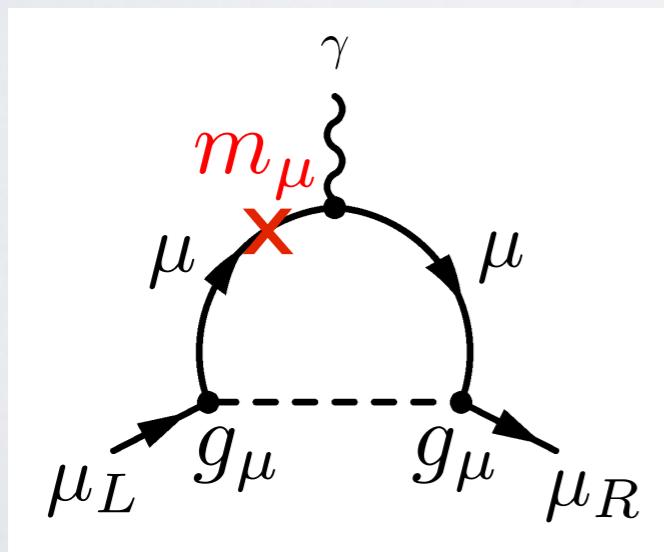
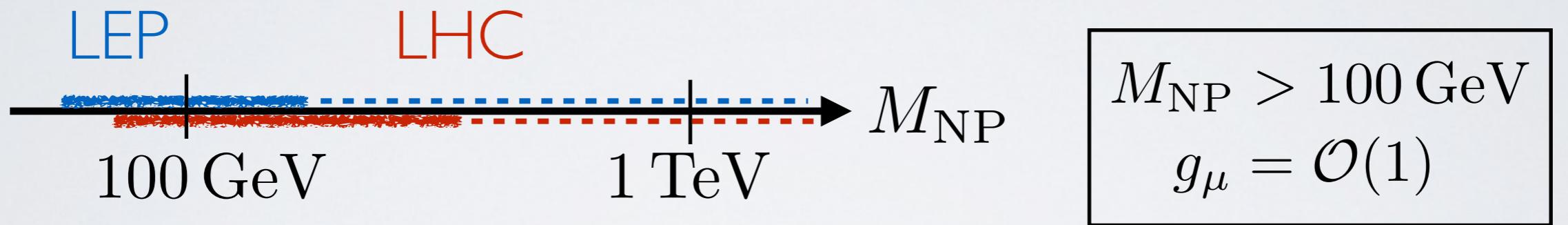
Discrepancy between standard model and experiment $\sim 3\sigma$.

NEW PHYSICS IN G-2

$$\Delta a_\mu \equiv a_\mu^{\text{exp}} - a_\mu^{\text{SM}} = (287 \pm 80) \times 10^{-11} \sim a_\mu^{\text{EW}}$$

[Jegerlehner, Nyffeler, arXiv:0902.3360]

Focus on mass range accessible in direct production at LHC:



Need chiral symmetry breaking:

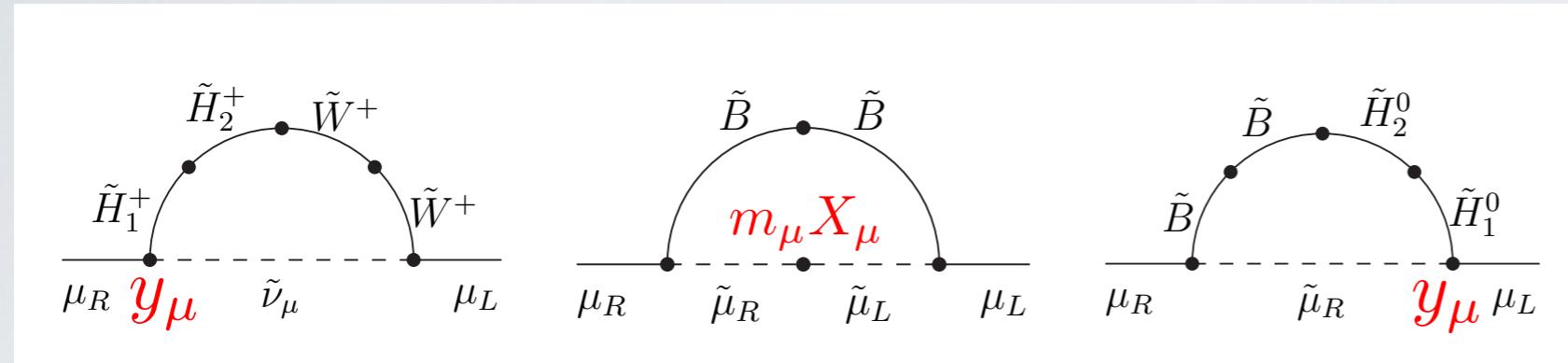
$$\Delta a_\mu^{\text{NP}} \sim g_\mu^2 \frac{m_\mu^2}{M_{\text{NP}}^2} \rightarrow \frac{a_\mu}{a_e} \sim \frac{m_\mu^2}{m_e^2} \sim 4 \times 10^4$$

[cf. Agrawal, Chacko, Verhaaren, arXiv:1402.7369]

PROMINENT EXAMPLES

Supersymmetric Standard Model

$\tan \beta$ enhancement from muon Yukawa coupling $y_\mu = \frac{m_\mu}{v} \tan \beta$



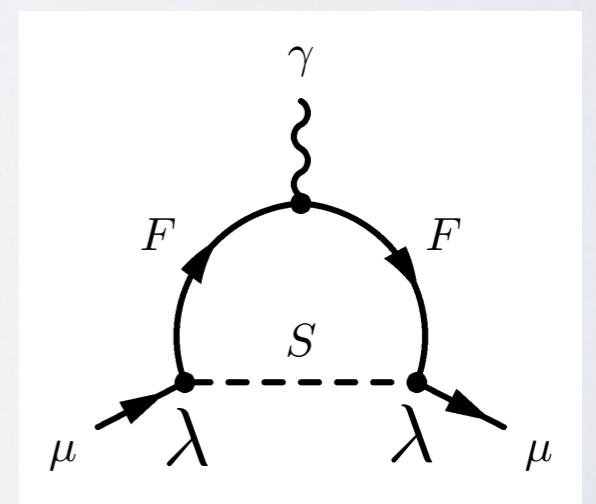
$$a_\mu \sim m_\mu^2 \tan \beta \mu M$$

[see, for instance, Stöckinger, J.Phys.G34 (2007) R45]

Light dark matter ($M_S \lesssim M_F \sim 1 \text{ GeV}$)

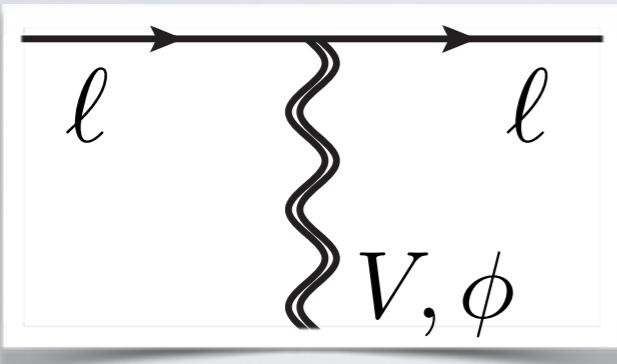
Δa_μ favors $M_F/\lambda \lesssim 100 \text{ GeV}$

Here: dark scalar plus fermion mediator

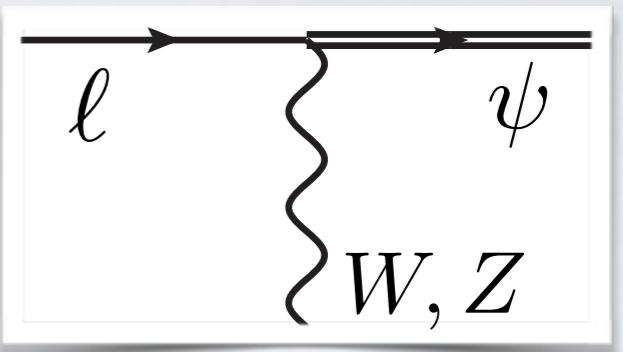


[e.g. Gnilenko 2001, Fayet 2007, Pospelov 2008, Agrawal, Chacko, Verhaaren, arXiv:1402.7369]

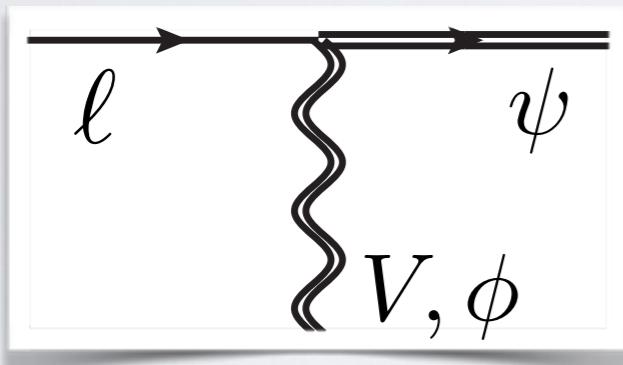
LEPTON INTERACTIONS WITH NEW FIELDS



new boson (scalar/vector)



new fermion



new boson and new fermion

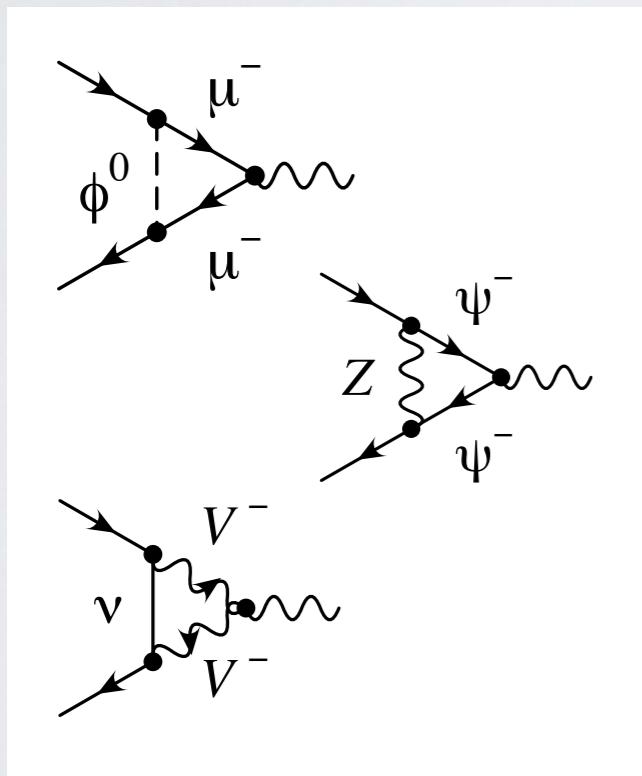
How can we test such g-2 candidates at colliders?

A MODEL-INDEPENDENT APPROACH

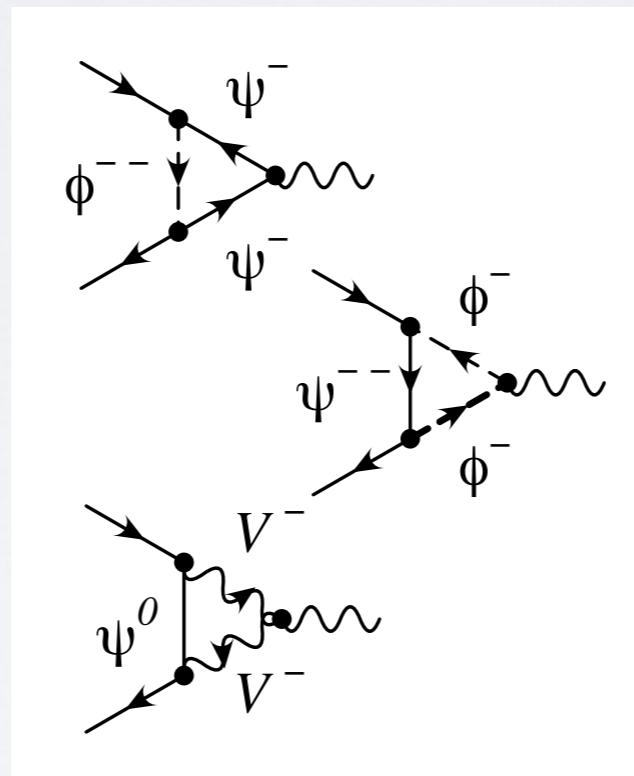
Consider new fields:

- spin 0, 1/2, 1 and integer electric charge
- weak singlets, doublets, triplets (all color singlets)

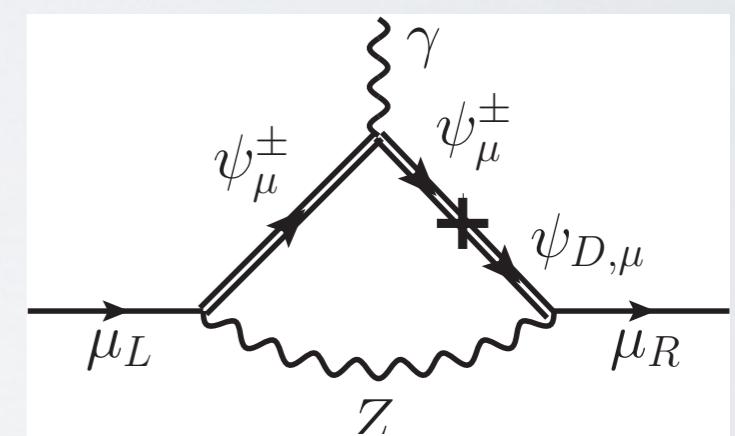
Three classes of one-loop contributions to a_μ :



One new field

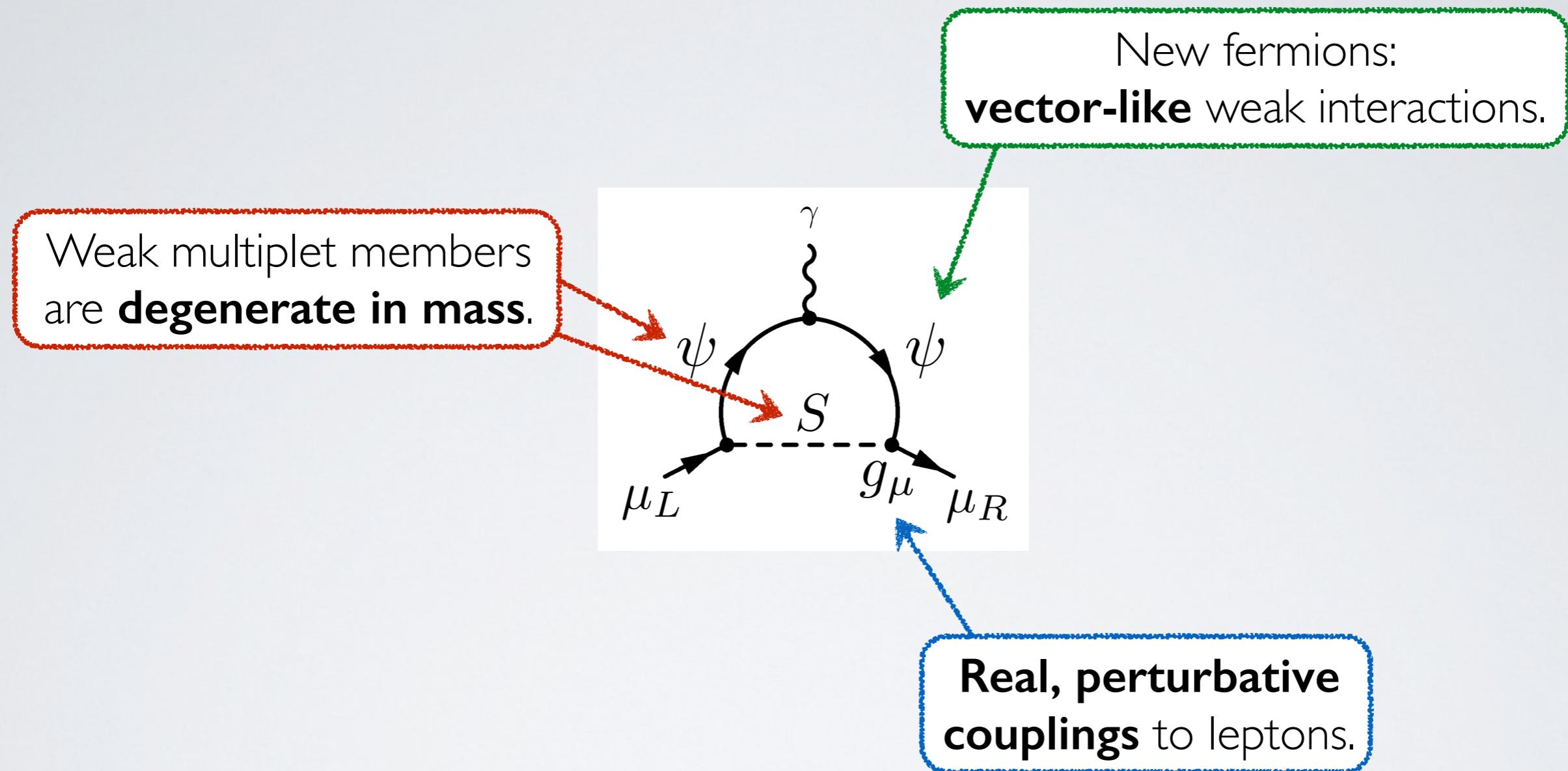


Two new fields



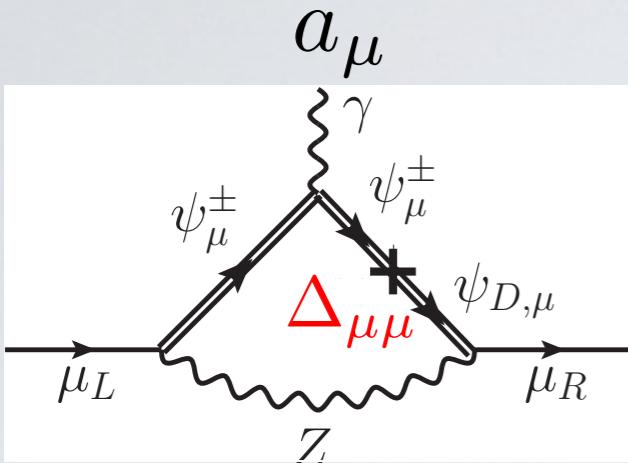
Two mixing fermions

GENERAL PARTICLE FEATURES

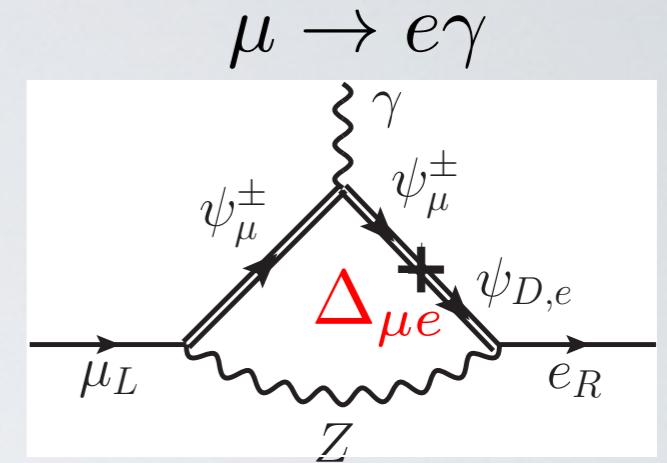


Assumption: Interactions are **minimally flavor-violating**.

MINIMAL FLAVOR VIOLATION



$$\mathcal{O}_{\mu\ell} = gy_\mu \Delta_{\mu\ell} H^\dagger \bar{\ell}_R \sigma^{\alpha\beta} A_{\alpha\beta} \mu_L$$



Strong experimental bound on flavor violation $\sim \Delta_{\mu e}$:

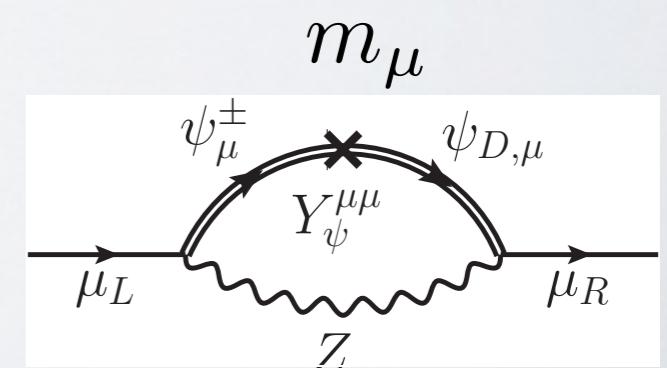
$$\mathcal{B}(\mu \rightarrow e\gamma) \approx 6.34 \times 10^{-7} \left(\frac{1 \text{TeV}^4}{\Lambda_{\text{FV}}^4} \right) |\Delta_{\mu e}|^2 < 5.7 \times 10^{-13} \quad [\text{MEG collaboration, 2013}]$$

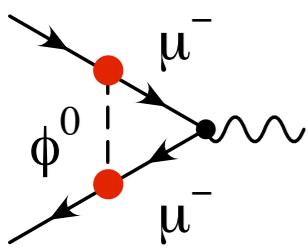
[Cirigliano et al., Nucl.Phys.B728 (2005) 121]

MFV implications for g-2 scenarios:

- Fermions in **fundamental** representation of $\mathcal{G}_F = \text{SU}(3)_L \times \text{SU}(3)_e$.
- Three **flavor copies** of new fermions.
- (At least) one Yukawa suppression, here: $\Delta_{\mu\mu} \sim y_\mu$.

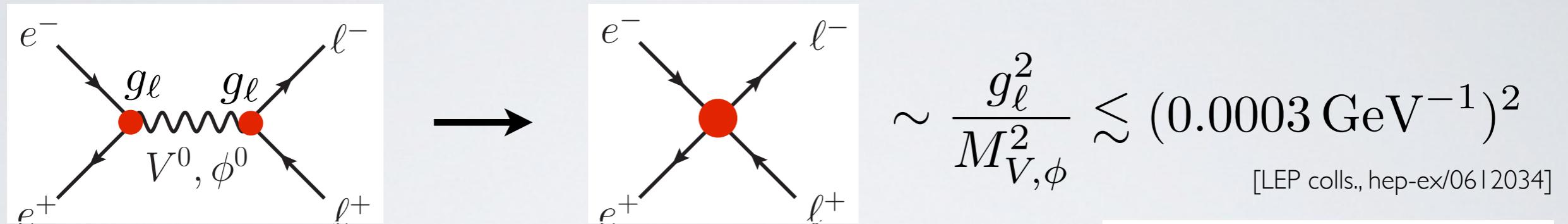
Muon mass protected: $m_\mu = y_\mu v (1 + \delta m_\mu)$





ONE NEW BOSON: LEP BOUNDS

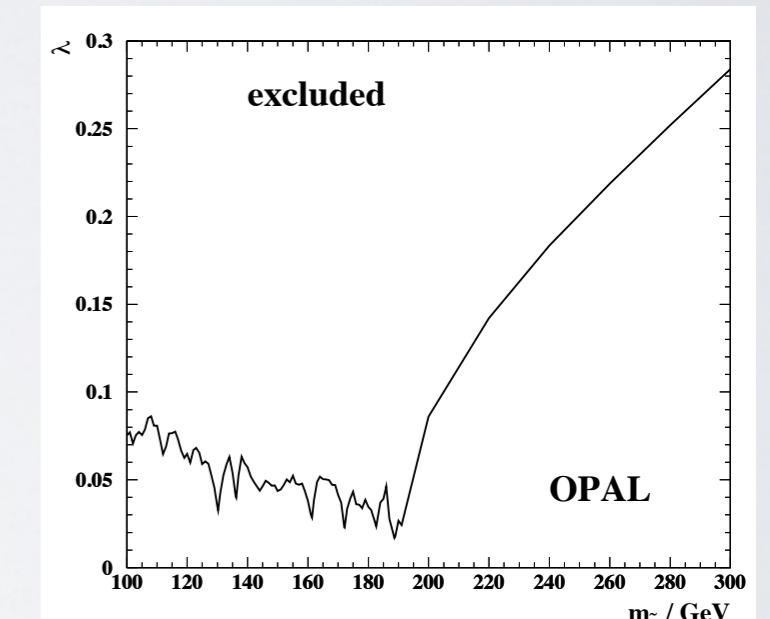
$M > \sqrt{s} \sim 200 \text{ GeV}$: 4-lepton contact interactions



$M < 200 \text{ GeV}$: neutralino resonance searches

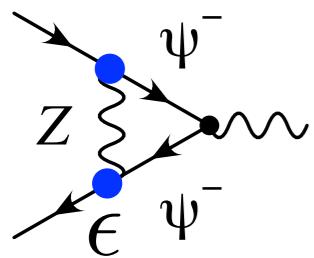
$$e^+ e^- \rightarrow (V^0, \phi^0) \gamma \rightarrow \ell^+ \ell^- \gamma$$

$$\rightarrow g_\ell / M_{V,\phi} \lesssim 0.0008 \text{ GeV}^{-1}$$



[OPAL coll., Eur.Phys.J. C13 (2000) 553]

Exclude neutral vector boson and scalar doublet
or scalar adjoint triplet as explanations of Δa_μ .



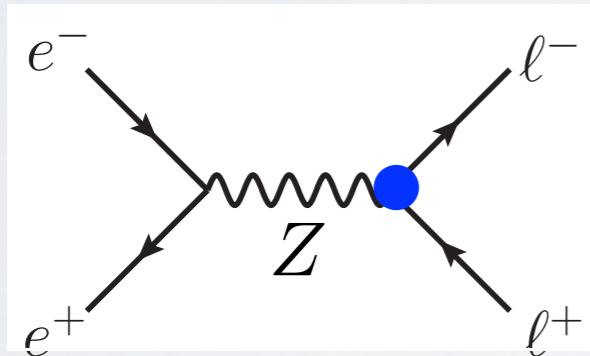
ONE NEW FERMION: LEP BOUNDS

New fermions mixing with SM leptons, $\epsilon = \frac{Y_\ell v}{M_\psi}$:

$$\mathcal{L} \supset -Y_\ell \overline{\psi_L} H \ell_R + \text{h.c.} \xrightarrow{\langle H \rangle} \epsilon \overline{\psi_R} \gamma^\mu \ell_R Z_\mu \quad (a_\mu)$$

and $(1 - \epsilon^2/2) \overline{\ell_R} \gamma^\mu \ell_R Z_\mu$ (LEP)

Global fit to LEP data: $|\epsilon| \lesssim 0.03$ [Aguila et al., PRD78 (2008) 013010]



$$\sim (1 - \epsilon^2/2)$$

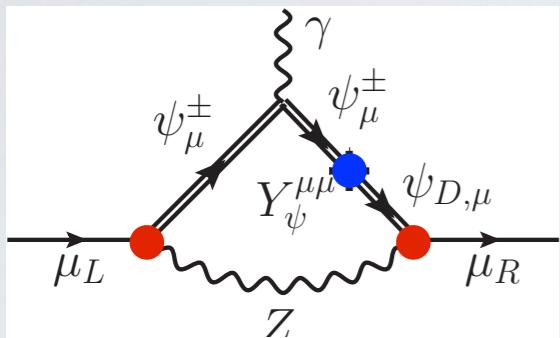
Excludes explanation of Δa_μ by a fermion doublet or charged singlet (neutral singlet and triplet yield $\delta a_\mu < 0$).

MIXING VECTOR FERMIONS

(Here: charged singlet ψ^\pm and doublet ψ_D)

Two viable scenarios of **minimal flavor violation**:

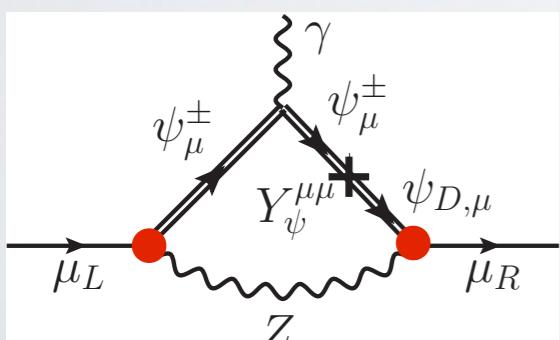
a) $\psi^\pm = (3, 1)$, $\psi_D = (1, 3)$



$$|Y_\psi| \lesssim 0.2, |\epsilon| \lesssim 0.03$$

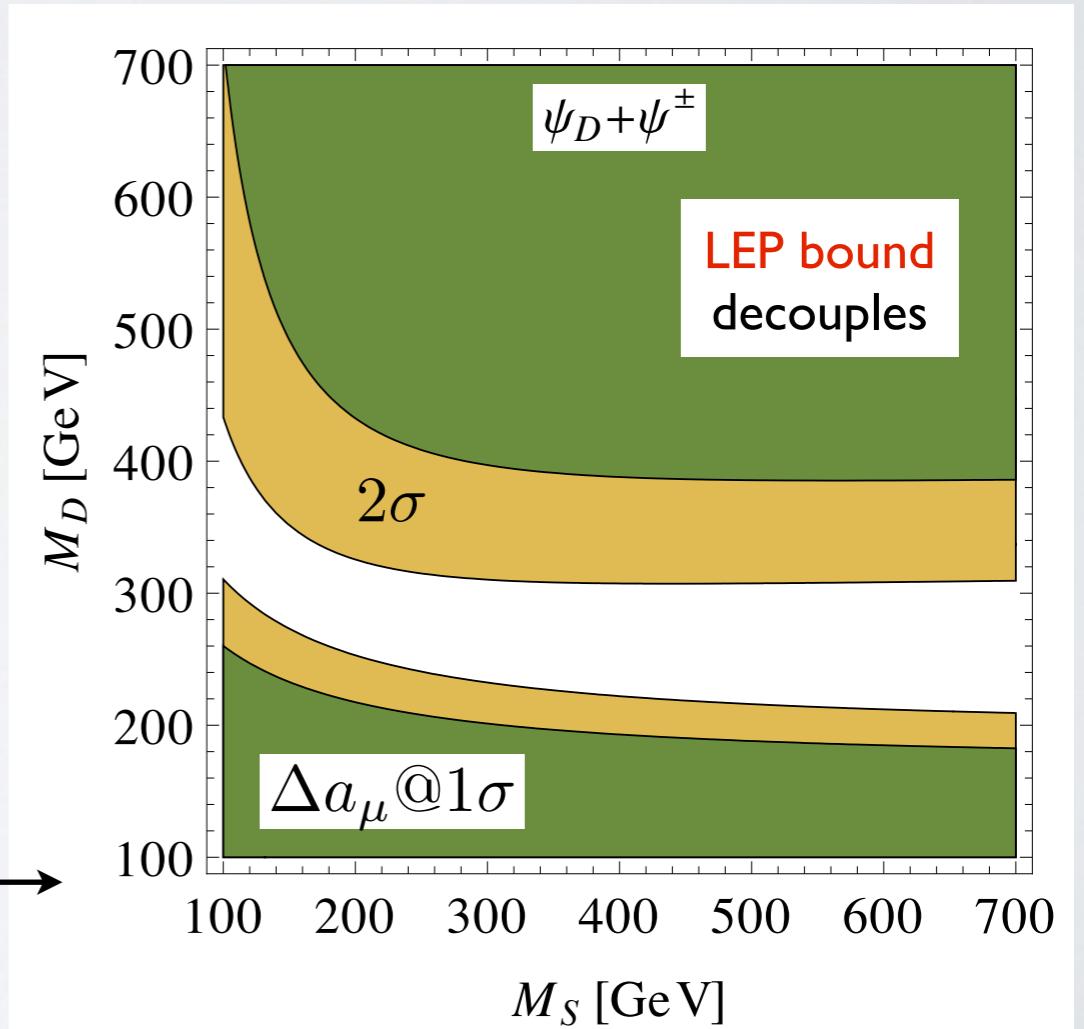
$$\rightarrow \delta a_\mu \lesssim \mathcal{O}(10^{-10})$$

b) $\psi^\pm, \psi_D = (1, 3)$ or $(3, 1)$



$$|Y_\psi| \lesssim \sqrt{4\pi}, |\epsilon| \lesssim 0.03$$

$$\delta a_\mu \sim m_\mu \cdot Y_\psi v \cdot \epsilon_1 \epsilon_2$$

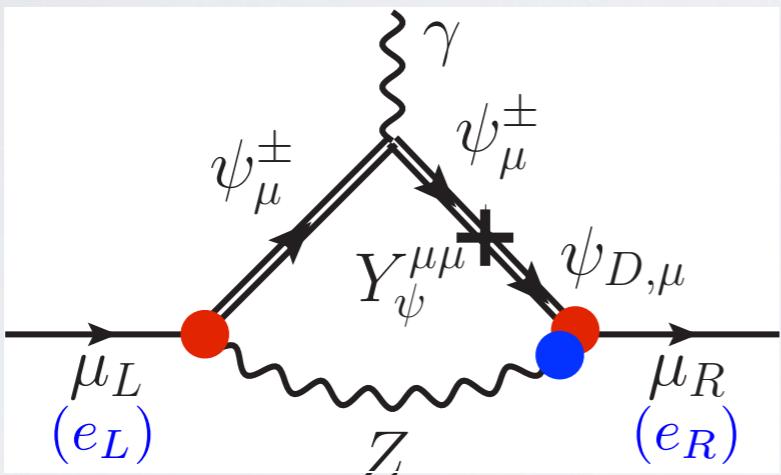


BOUNDS FROM ELECTRON MOMENT

(Here: charged singlet ψ^\pm and doublet ψ_D)

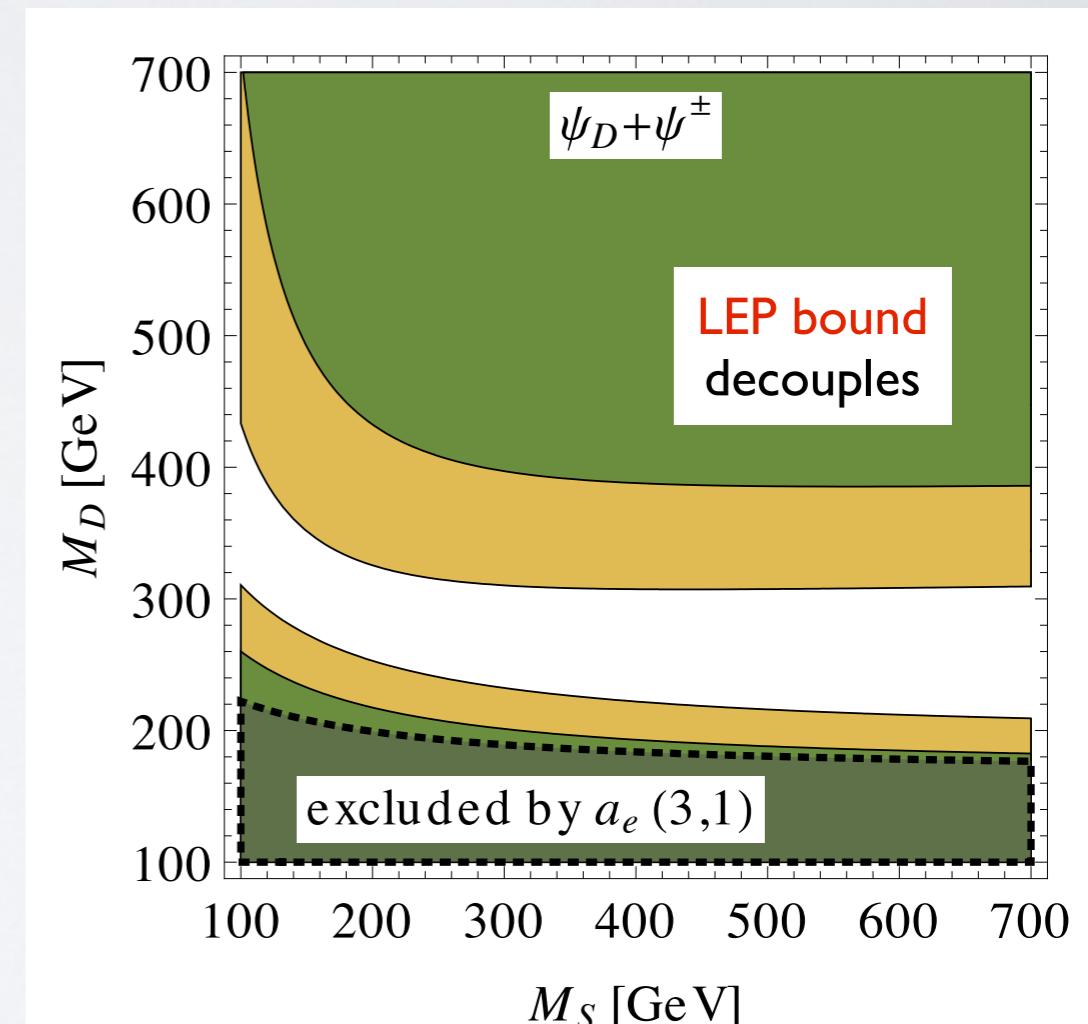
Electron g-2: $a_e^{\text{exp}} - a_e^{\text{SM}} = (-1.06 \pm 0.82) \times 10^{-12}$ [Aoyama et al., PRL 109 (2012) 111807]

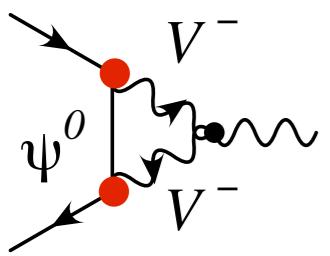
Consider $\psi^\pm = \psi_D = (3, 1)$:



$$\delta a_e = \frac{m_e^2}{m_\mu^2} \cdot \max \left\{ \frac{\sqrt{4\pi} y_\mu / y_\tau}{0.03 M_D / v}, 1 \right\} \delta a_\mu$$

$= \mathcal{O}(10^{-12})$ for $M_D \lesssim v$





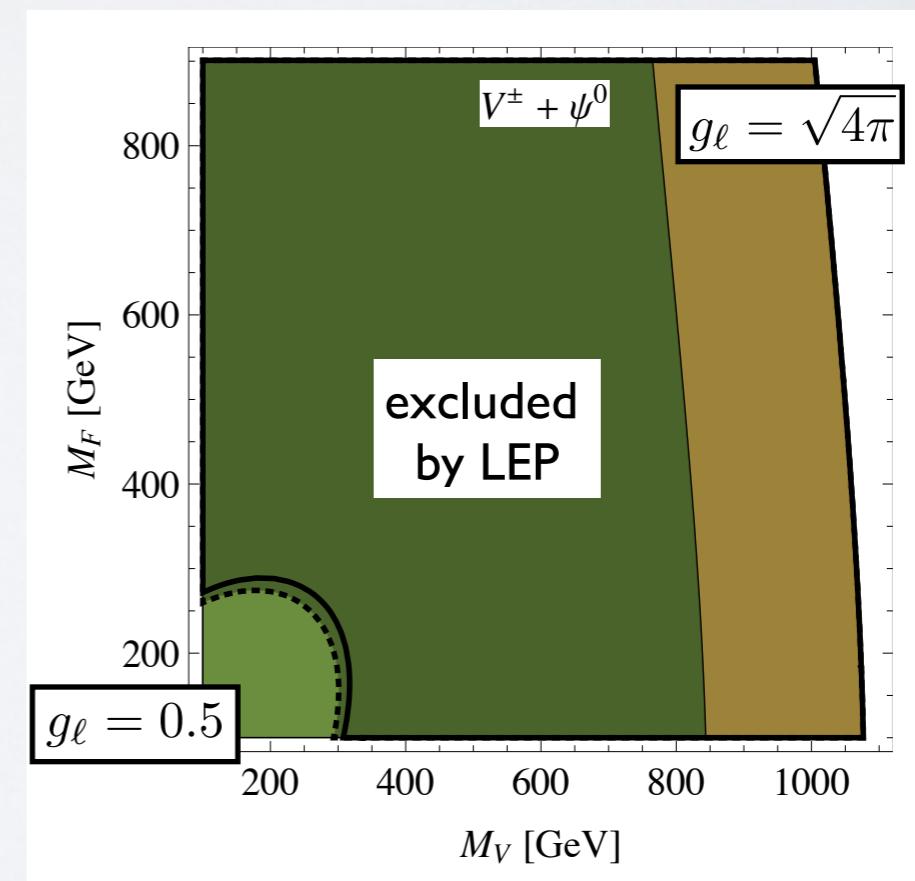
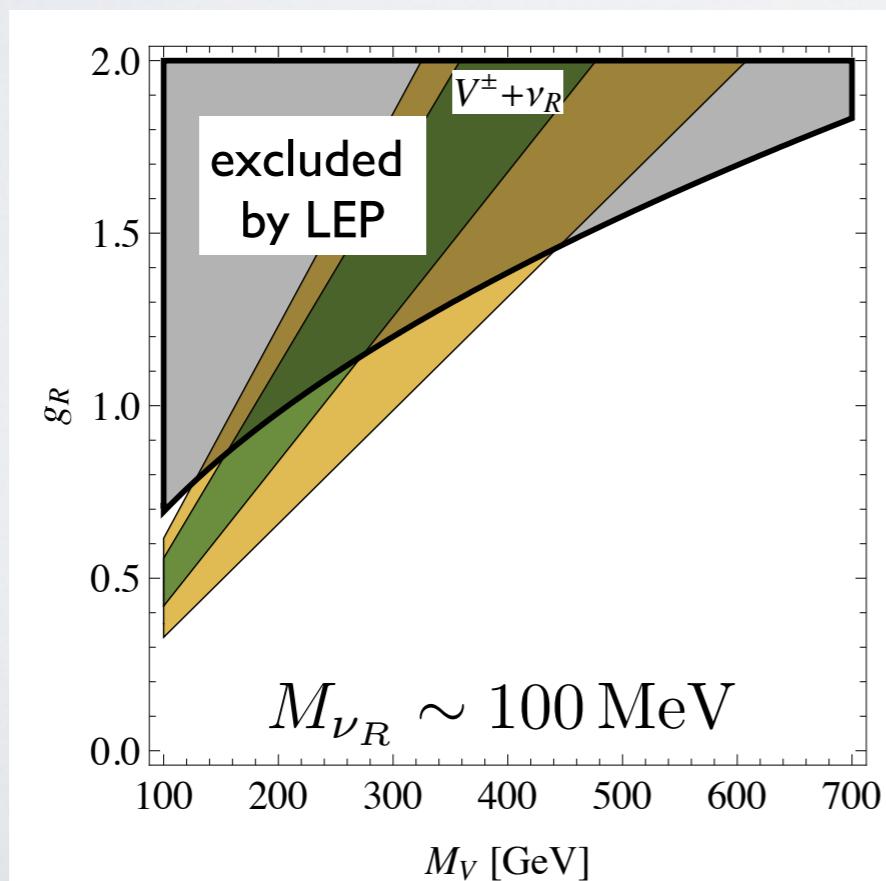
TWO NEW FIELDS: LEP BOUNDS

Indirect bound on SM-NP-NP couplings at one-loop level:

$$\sim \frac{g_\ell^4}{16\pi^2 M_V^2} \lesssim (0.0003 \text{ GeV}^{-1})^2$$

[LEP colls., hep-ex/0612034]

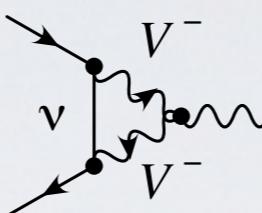
LEP constraints on δa_μ in regions of strong couplings:



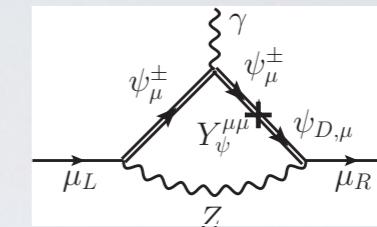
VIABLE G-2 CANDIDATES AFTER LEP

One new field

- charged vector singlet V^\pm

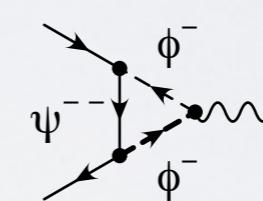
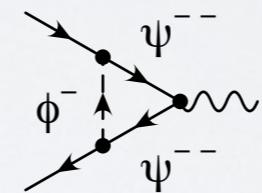
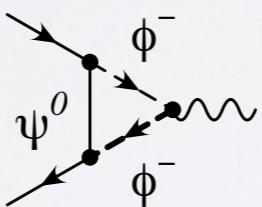
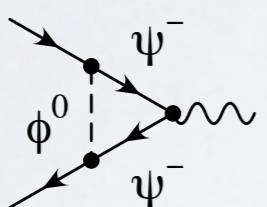


Mixing fermions

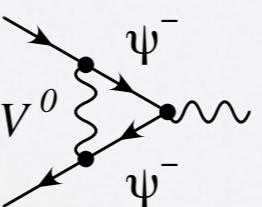
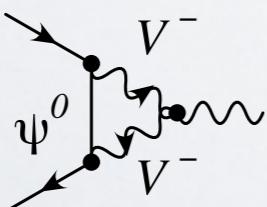


Two new fields

- $\phi^0 + \psi^\pm$: neutral scalar + charged fermion singlet
- $\phi_D + \psi_{A,T}$: scalar doublet & fermion triplets ($Y=0$ and $Y=-1$)
- $\phi_A + \psi_T$: scalar adjoint triplet & fermion triplet ($Y=-1$)



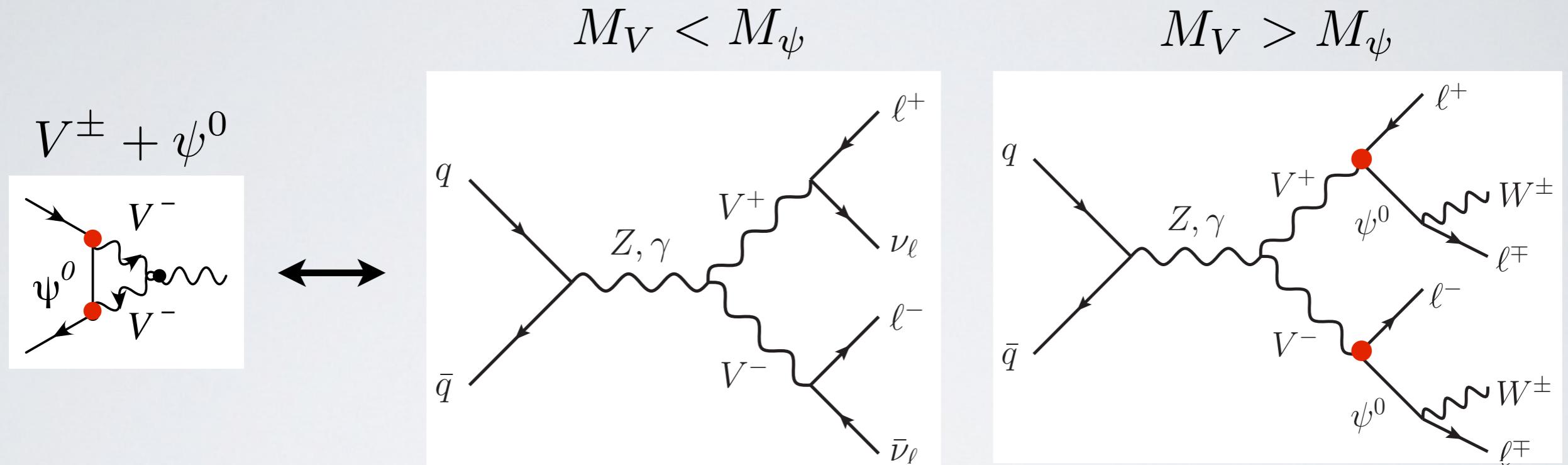
- $V^\pm + \psi^0$: charged vector singlet & neutral fermion
- $V_A + \psi_D$: vector adjoint triplet & fermion doublet



One-loop LEP bounds optional;
may be relaxed by tree-level NP.

DIRECT TESTS AT THE LHC

(Here: charged vector singlet & neutral fermion singlet)



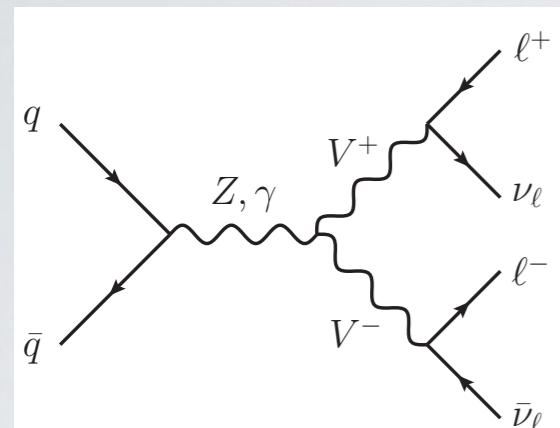
Assumptions

- Drell-Yan pair production (model-independent).
- Two new fields: constrain the lighter one (no cascades).
- If this is a SM singlet, look for cascade decays.
- Decay dominantly (and flavor-universally) into leptons.

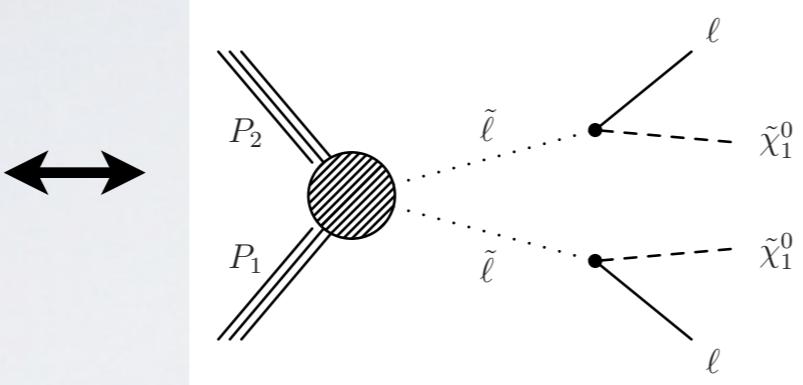
RECAST SEARCHES BY ATLAS & CMS

(charged vector singlet & neutral fermion singlet ctd.)

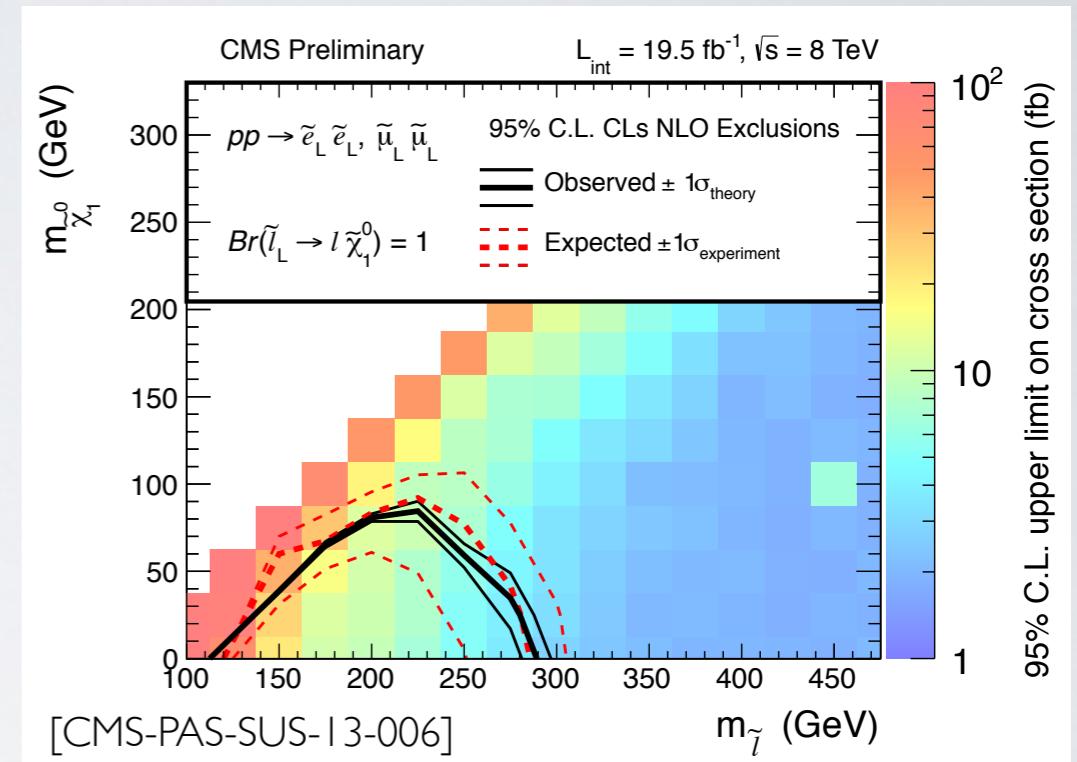
$M_V < M_\psi$:



slepton production



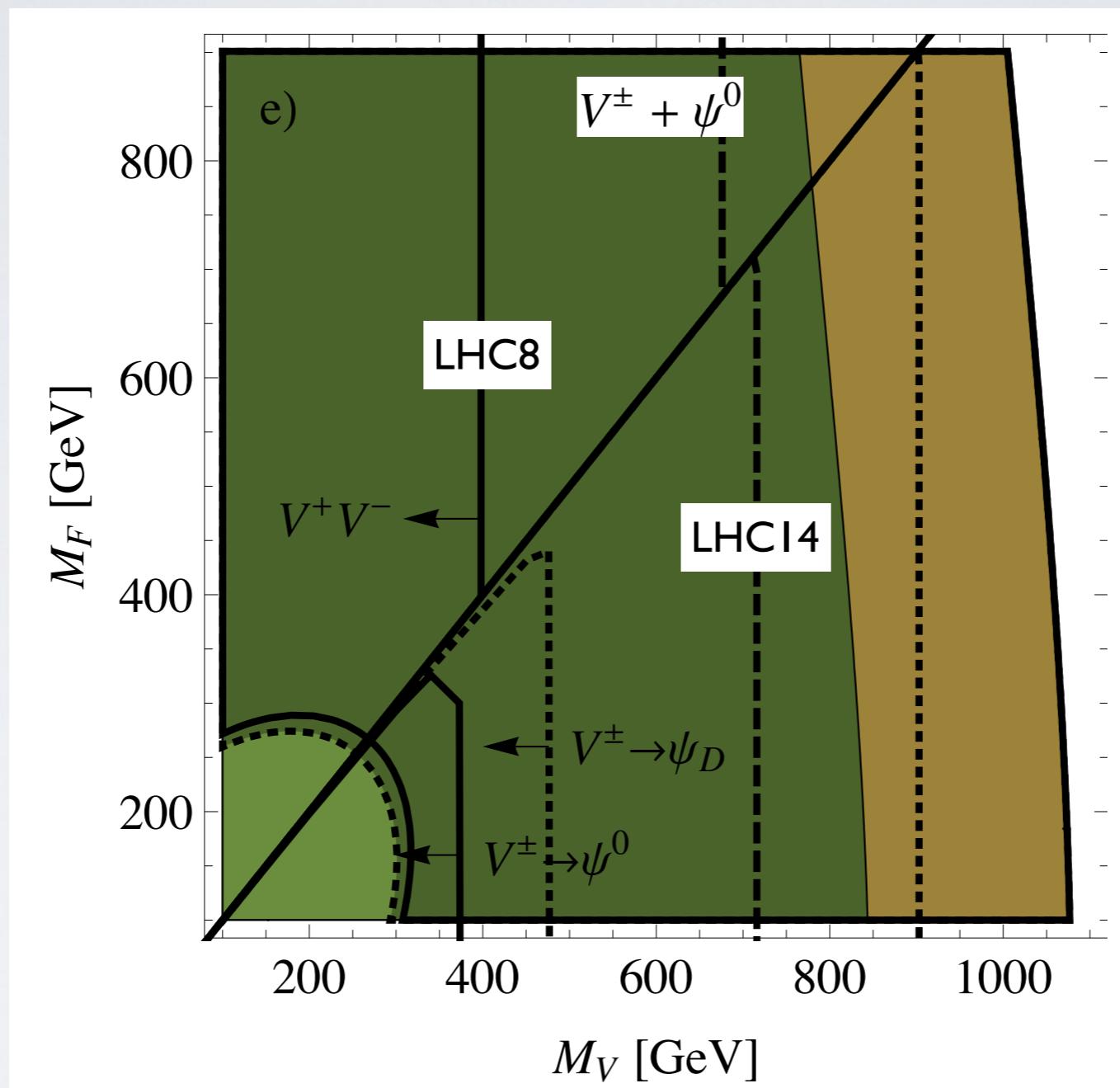
$\rightarrow M_V > 398 \text{ GeV}$



$M_V > M_\psi$: multi-lepton search [ATLAS-CONF-2013-019]

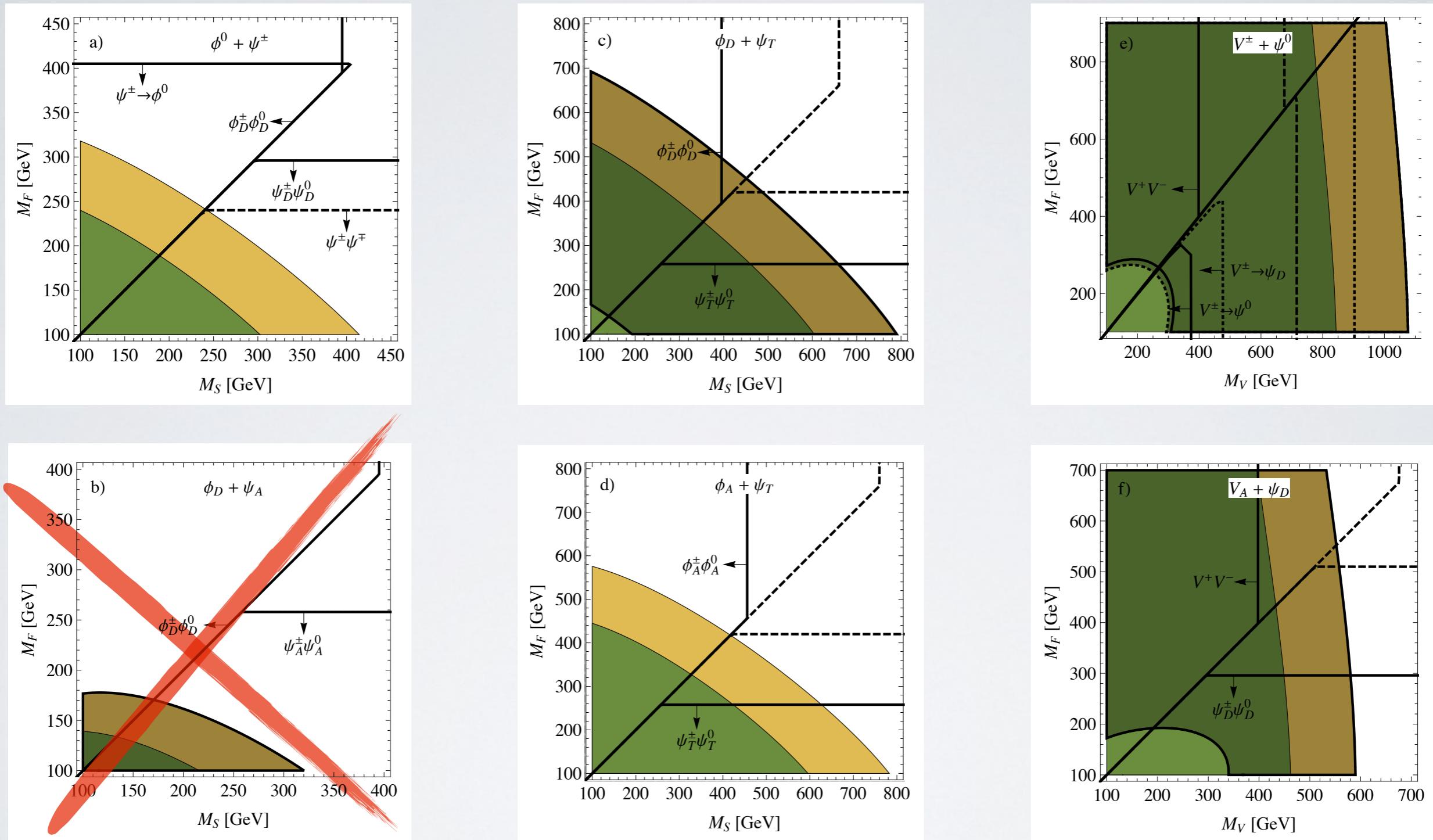
- ψ^0 is weak singlet: $\mathcal{B}(\psi^0 \rightarrow W^\pm \ell^\pm) = 1/2$ $M_V > 373 \text{ GeV}$
- ψ^0 part of weak doublet: $\mathcal{B}(\psi^0 \rightarrow W^\pm \ell^\pm) = 1$ $M_V > 476 \text{ GeV}$

LHC CONSTRAINTS ON G-2



LHC14: rescaled event yield $N_{14}(300 \text{ fb}^{-1}) = N_8(20 \text{ fb}^{-1}) \frac{\sigma(14 \text{ TeV})}{\sigma(8 \text{ TeV})}$

G-2 CANDIDATES AFTER LHC8



Mixing fermions for g-2 cannot be excluded at the LHC.

PROSPECTS: ORIGIN OF G-2 ANOMALY?

- Current g-2 uncertainties:

$$\delta a_{\mu}^{\text{exp}} \approx \pm 63 \times 10^{-11}$$

$$\delta a_{\mu}^{\text{SM}} \approx \pm 49 \times 10^{-11}$$

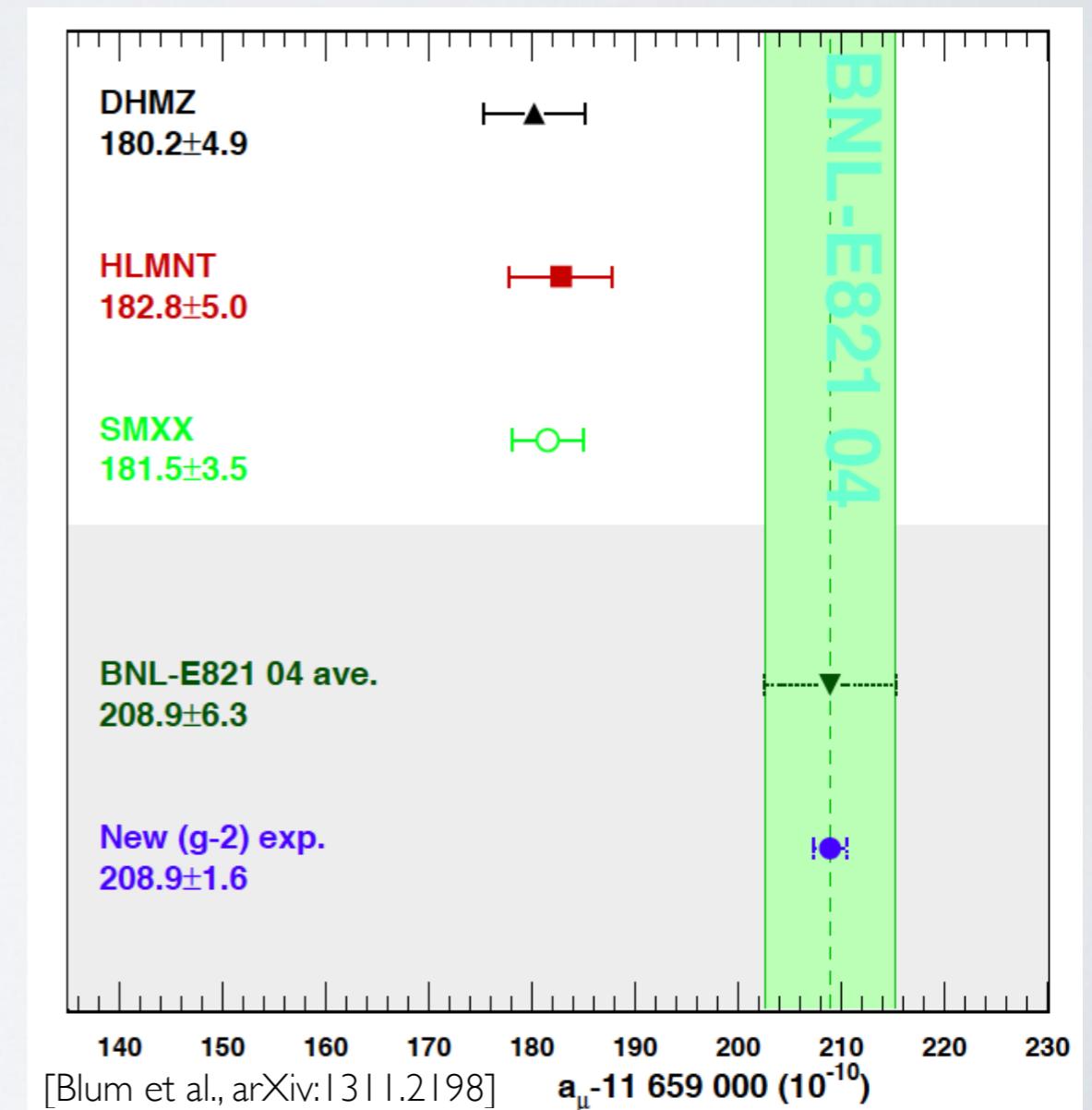
- Fermilab's E989 and J-PARC:

$$\delta a_{\mu}^{\text{exp}} \approx \pm 16 \times 10^{-11}$$

- Expected reduced SM error:

$$\delta a_{\mu}^{\text{SM}} \approx \pm 35 \times 10^{-11}$$

Total uncertainty may be reduced by a factor of 2.



SUMMARY

Combined with LEP bounds,
all of our simplified models for g-2 but mixing fermions
can be **conclusively tested at LHC during run II.**

In case Fermilab and J-PARC find indirect evidence of new
particles in g-2, the **LHC might provide**
the only (?) way to discover them directly.
